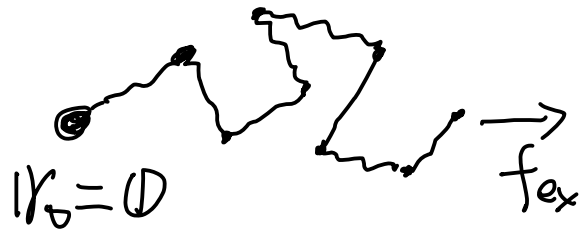


系統計力学A IX

講義 X 毛

2023 / 11 / 28

§ 今日の目標



・ (3次元空間中の) 1次元金鎖

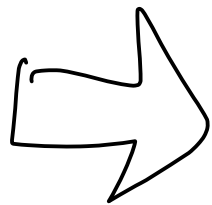
・ $\Gamma = (r_1, r_2, \dots, r_N, p_1, \dots, p_N)$

$r = (x, y, z)$

$$H(\Gamma; f_{ex}) = H_0(\Gamma) - \underbrace{f_{ex} x_N}_{\text{一定外力}}$$

例: $H_0(\Gamma) = \sum_{i=1}^N \frac{(p_i)^2}{2m} + \frac{k}{2} \sum_{i=1}^N (|r_i - r_{i-1}| - a)^2$

自由に回転できる "かたい" はね
 $\beta k a^2 \gg 1$



統計力学へ

§ カノニカル分布

$$\mathcal{P}_{\beta, f_{\text{ex}}}^c(\Gamma) = \frac{1}{Z(\beta, f_{\text{ex}}, N)} e^{-\beta \underbrace{(H_0(\Gamma) - f_{\text{ex}} x_N)}_{H(\Gamma; f_{\text{ex}})}}$$

- $X \equiv \langle x_N \rangle_{\beta, f_{\text{ex}}}^c = k_B T \frac{\partial}{\partial f_{\text{ex}}} \log Z(\beta, f_{\text{ex}}, N)$

- $E \equiv \langle H \rangle_{\beta, f_{\text{ex}}}^c = - \frac{\partial}{\partial \beta} \log Z(\beta, f_{\text{ex}}, N)$

$$Z(\beta, f_{\text{ex}}, N) = \int d\Gamma e^{-\beta H(\Gamma; f_{\text{ex}})}$$

$$= e^{-\beta (E_* - T S(E_*, f_{\text{ex}}, N))}$$

$$E_*(T, f_{\text{ex}}, N): \left(\frac{\partial S}{\partial E} \right)_{E_*, f_{\text{ex}}, N} = \frac{1}{T} \rightarrow$$

$$\underline{E_* = E}$$

$$S(E, f_{\text{ex}}, N) = k_B \log \Sigma(E, f_{\text{ex}}, N) + o(N) \\ (= k_B \log \Omega(E, f_{\text{ex}}, N))$$

N! くらい
→

? 章のギョウ
と似て

& 3冊目
page 12
を参照

§ 自由エネルギー - $\tilde{H}$

$$\tilde{H}(T, f_{ex}, N) \equiv -k_B T \log Z(T, f_{ex}, N) \quad \beta = \frac{1}{k_B T}$$

$$\left\{ \begin{array}{l} X = - \frac{\partial \tilde{H}}{\partial f_{ex}} \\ E = \frac{\partial}{\partial \beta} (\beta \tilde{H}) = \tilde{H} - T \frac{\partial \tilde{H}}{\partial T} \\ \tilde{H} = E - TS \end{array} \right.$$
$$\Rightarrow S = - \frac{\partial \tilde{H}}{\partial T}$$

$$\underline{d\tilde{H} = -SdT - Xdf_{ex}}$$

熱力学関係式

§ 自由エネルギー-H: 熱力学

$$dH = -SdT - f dX \quad (1): f: \text{復元力}$$

$$T, S \quad d\tilde{H} = -SdT - X df_{ex} \quad (2)$$

$$= -SdT + X df \quad \text{??? } f_{ex} + f = 0$$

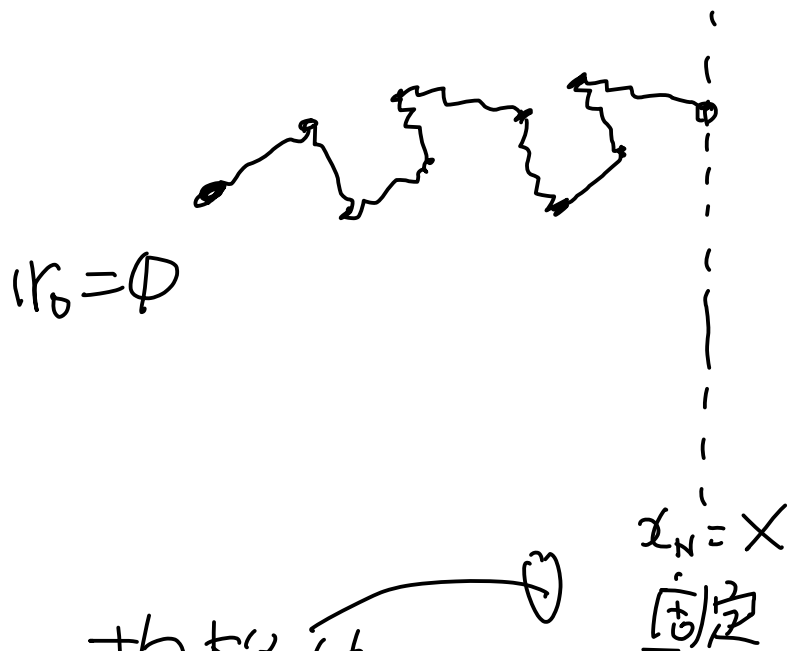
$$\tilde{H} = \underbrace{H - f_{ex} X}_{\text{外力のポテンシャル}}$$

$$\begin{aligned} (1) \Rightarrow d\tilde{H} &= -SdT - f dX - f_{ex} dX - X df_{ex} \\ &= -SdT - X df_{ex} \quad (2) \end{aligned}$$

$\tilde{H}$: 外力のポテンシャルを加えた自由エネルギー

⇒ H の 統計力学的導出

§ 設定



- 拘束条件
- ポテンシャルを用いた
極限操作を以て
表現可能
(“壁”と同じ)

✓ 力 = カル分布

$$\rho_{\beta, X, N}^C(\Gamma) = \frac{1}{Z_{\text{II}}(\beta, X, N)} e^{-\beta H_0(\Gamma)} \cdot \underbrace{\delta(x_N - X)}_{\text{拘束条件}}$$

✓ 復元力

$$f = - \left\langle \frac{\partial H_0(\Gamma)}{\partial x_N} \right\rangle_{\beta, X, N}^C$$

$$Z_{\text{II}}(\beta, X, N) = \int d\Gamma e^{-\beta H_0(\Gamma)} \delta(x_N - X)$$

§ 復元力の公式

$$f = -\frac{1}{Z_H} \int dP \frac{\partial H_0}{\partial x_N} e^{-\beta H_0(P)} \delta(x_N - x)$$

$$= +\frac{k_B T}{Z_H} \int dP \frac{\partial}{\partial x_N} \left(e^{-\beta H_0(P)} \right) \cdot \delta(x_N - x)$$

$$= \frac{k_B T}{Z_H} \int dP e^{-\beta H_0(P)} \frac{\partial}{\partial x} \delta(x_N - x)$$

$$\begin{aligned} \frac{\partial}{\partial x_N} \delta(x_N - x) \\ = -\frac{\partial}{\partial x} \delta(x_N - x) \end{aligned}$$

$$= k_B T \frac{1}{Z_H} \frac{\partial}{\partial x} Z_H$$

$$= k_B T \frac{\partial}{\partial x} \log Z_H //$$

しかし、 Z_H の計算困難
(直接積分できない)

§ 熱力学関係式

$$\bar{F}(T, X, N) \equiv -k_B T \log Z_{\square}(T, X, N)$$

$$f = - \frac{\partial \bar{F}}{\partial X} \quad //$$

また、
 (6) の計算を繰り返して

$$\left\{ \begin{array}{l} E_0 = \langle H_0 \rangle_{\beta, X, N}^c = - \frac{\partial}{\partial \beta} \log Z_{\square}(\beta, X, N) \\ Z_{\square} = e^{-\beta (E_0 - TS(E_0, X, N))} \cdot \omega(N) \end{array} \right.$$

$$\Rightarrow S = - \frac{\partial \bar{F}}{\partial T} \quad ; \quad \bar{F} = E_0 - TS$$

$$\Rightarrow \underline{d\bar{F} = -SdT - f dX}$$

$$\begin{aligned}
 \psi = 3\tilde{z}, \quad dE_0 &= d\tilde{H} + TdS + SdT \\
 &= -TdS - fdx + TdS \\
 &= TdS - fdx \quad // \quad \star
 \end{aligned}$$

$$\begin{aligned}
 dE &= d\tilde{H} + TdS + SdT \\
 &= TdS - x d\tau_x \quad // \quad \star\star
 \end{aligned}$$

10x

$$E_0 = \langle H_0 \rangle_{\beta, x, N}^c : \text{内部エネルギー} - U$$

$$E = \langle H \rangle_{\beta, \tau_x, N}^c : \text{エンタルピー} - H$$

E は 内部エネルギー と呼ぶのは間違いないが、
2つの関係式 $\star, \star\star$ は区別すべき

$\mathcal{Z}$ と $\mathcal{Z}_H$ の関係

$$\mathcal{Z}(T, f_{ex}, N) = \int dP e^{-\beta H_0(P) + \beta f_{ex} x_N} \quad \leftarrow \text{計算機}$$

$$= \int dP \underbrace{\int dx \delta(x_N - x)}_{=1} e^{-\beta H_0(P) + \beta f_{ex} x_N}$$

$$= \int dx e^{\beta f_{ex} x} \int dP \delta(x_N - x) e^{-\beta H_0(P)}$$

$$= \int dx e^{\beta f_{ex} x - \beta F(T, x, N)}$$

$$= e^{\beta (f_{ex} x_* - F(T, x_*, N))}$$

$$= e^{\beta f_{ex} x_*} \mathcal{Z}_H(T, x_*, N)$$

追加 page 12
 を参照

$$\left. \frac{\partial}{\partial x} (f_{ex} x - F(T, x, N)) \right|_{x_*} = 0$$

$$\Leftrightarrow f_{ex} + \left. \frac{\partial F}{\partial x} \right|_{x_*} = 0$$

$$\Leftrightarrow f_{ex} + f = 0$$

$$\Leftrightarrow \tilde{H}(T, f_{ex}, N) = F(T, x_*, N) - x_* f_{ex}$$

§ まとめ

- $$P_{\beta, \text{tex}, N}^C(\Gamma) = \frac{1}{Z} e^{-\beta H_0(\Gamma) + \beta \text{tex} x_N}$$

↑ T-P 分岐とも
よばれる

- $$P_{\beta, x, N}^C(\Gamma) = \frac{1}{Z_{\mathbb{I}}} e^{-\beta H_0(\Gamma)} \delta(x_N - x)$$

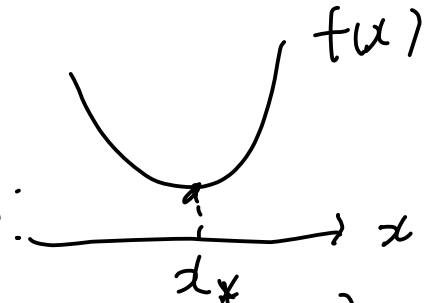
↑ P-ハンブル
の等価性

指数関数の肩に $\frac{2}{\beta}$ 数をあげ"子=2"

計算できるようになる。

(三ノロカニカシ ~~カニカシ~~ カニカシ も同様)

§ 鞍点法

積分 $I = \int dx e^{-Nf(x)}$ の評価: 

$$\begin{aligned} I &= \int dx e^{-N \left[f(x_*) + \frac{1}{2} f''(x_*) (x-x_*)^2 + \dots \right]} \\ &= e^{-Nf(x_*)} \underbrace{\int dx e^{-N \left[\frac{1}{2} f''(x_*) (x-x_*)^2 + \dots \right]}}_{\text{Gauss 積分} + \dots} \rightarrow O(\sqrt{N}) \\ &= e^{-Nf(x_*) + o(N)} \end{aligned}$$

↖ $f(x)$ の最小値と交点か?

例

$$\begin{aligned} &\int dE e^{\frac{1}{k_B} S(E) - \beta E} \\ &= N \int du e^{N \left(\frac{1}{k_B} \bar{s}(u) - \frac{1}{k_B T} u \right)} \\ &E = Nu; S(E) = N \bar{s}\left(\frac{E}{N}\right) \end{aligned}$$

$$\begin{aligned} &\int dx e^{\beta f_{ex} x - \beta F(T, x)} \\ &= N \int dx e^{N \left(\beta f_{ex} x - \beta \phi(T, \frac{x}{N}) \right)} \\ &x = Nx; F(T, x) = N \phi(T, \frac{x}{N}) \end{aligned}$$