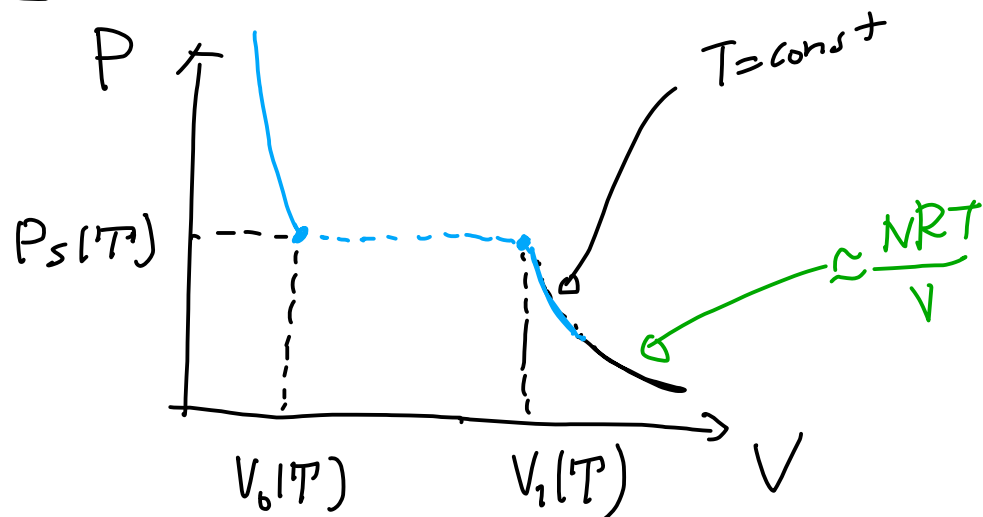
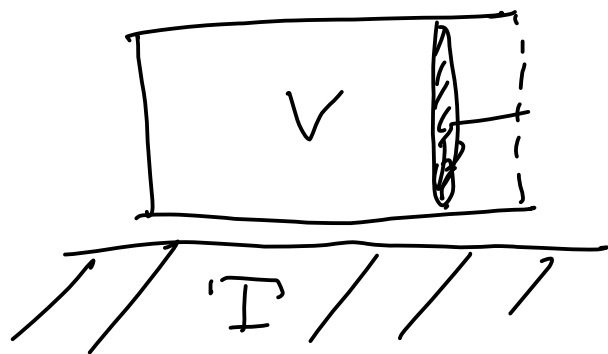


熱力学講義X

20/07/15

§ 設定



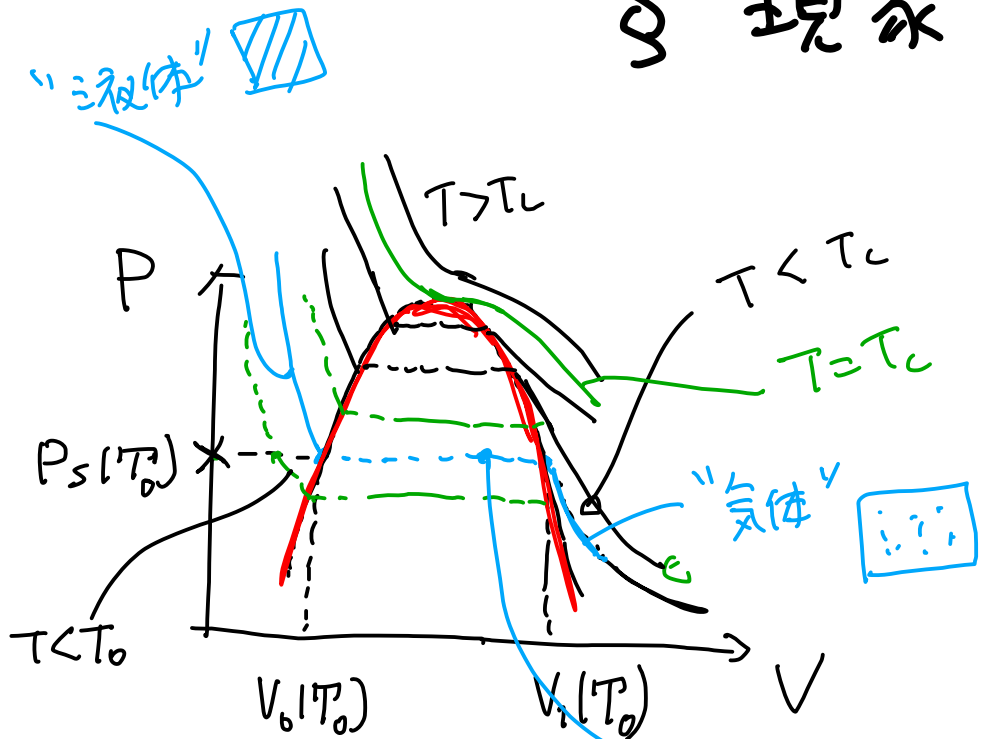
• $V \nearrow$

$$P \approx \frac{NRT}{V}$$

$P_s(T)$: 飽和圧力

この現象を熱力学
として議論する。

§ 現象

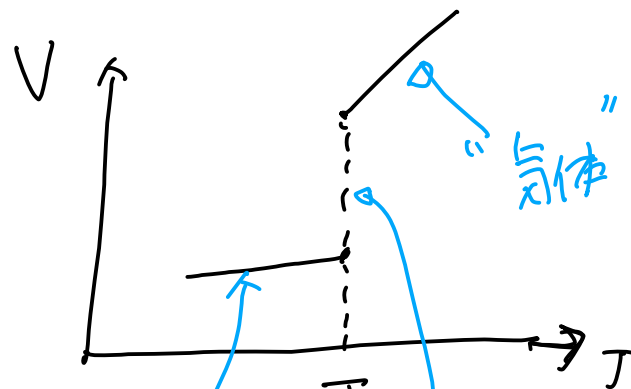
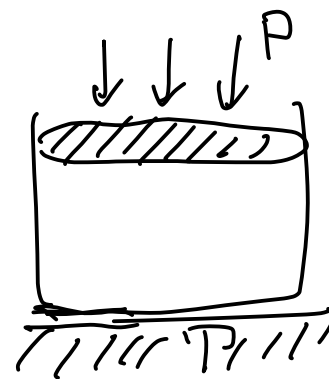


T_c : 臨界溫度

氣液共存



$P = \text{const}$
 (e.g. $P = P_s(T_0)$)



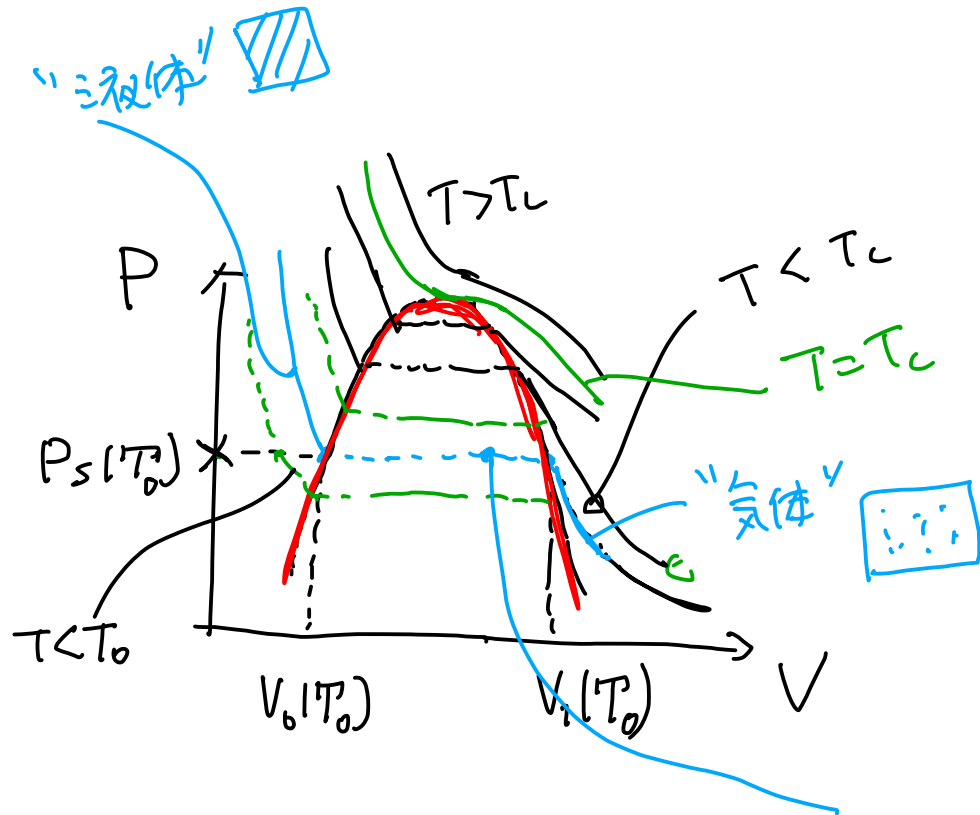
$P = P_s(T_0)$

$(\Rightarrow T_0 = T_0(P))$
 轉移溫度

氣液轉移

~ Intermission ~

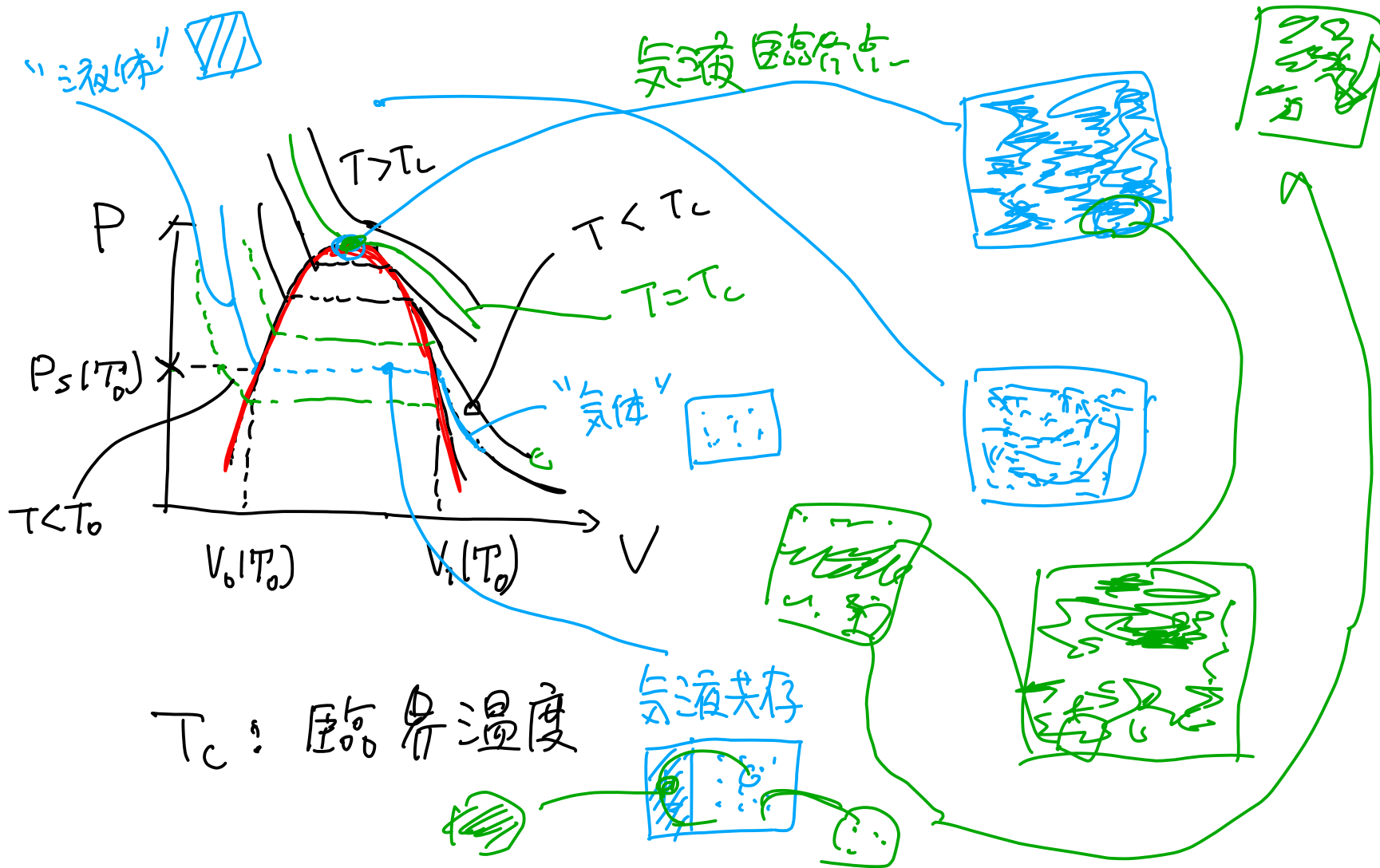
§ 気体と液体の区別？




T_c : 臨界温度

気液共存

臨界点の様子

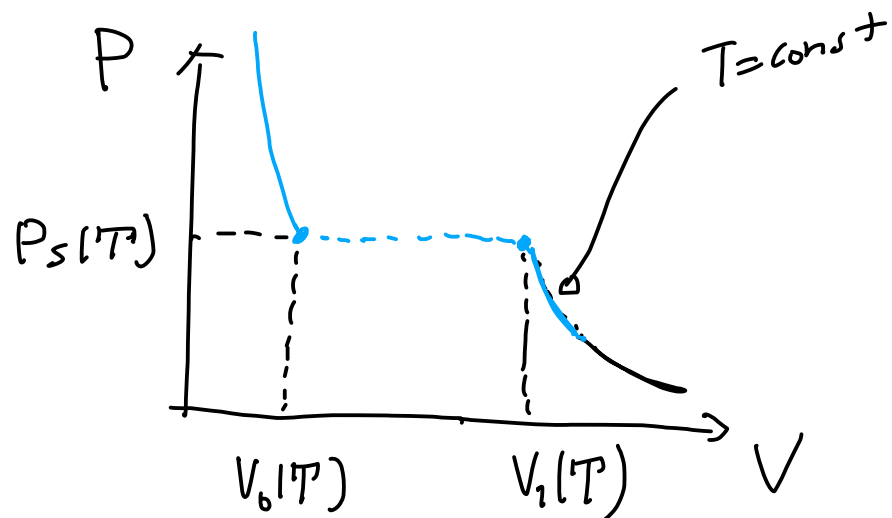


• 自己相似形

• フラクタル


~ Intermission ~

§ クラペイロンの式



例:

水 \rightarrow 水蒸気

1気圧 \rightarrow 373 K

1500 m の山
1気圧 - 150 hPa \rightarrow 368 K

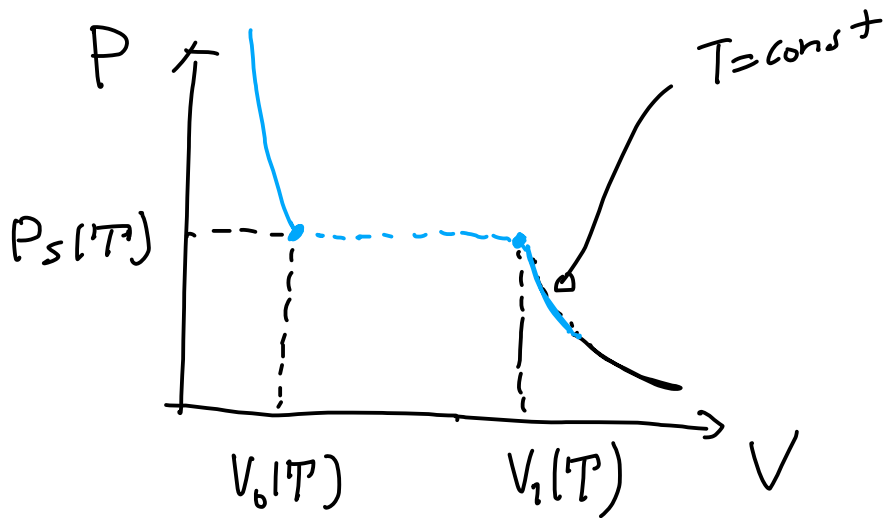
$$T_0 = T_0(P)$$

$$\Leftrightarrow P_s = P_s(T)$$

L: 気化熱

$$\left(\frac{dT_0(P)}{dP} \right)^{-1} = \frac{dP_s(T)}{dT} = \frac{L}{T(V_1 - V_0)}$$

§ 気化熱



$$L = Q[(T, V_0(T)) \xrightarrow{ig} (T, V_1(T))] \\ = T [S(T, V_1(T)) - S(T, V_0(T))]$$

分子間エネルギー ↑
 分子の運動エネルギー →
 “熱” を吸収

$$d\bar{H} = -SdT - PdV$$

§ 導出

$$\bar{H}(T, V_1(T)) - \bar{H}(T, V_0(T)) = -P_S(T)(V_1(T) - V_0(T))$$

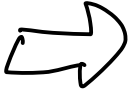
T での微分

$$\left(\frac{\partial \bar{H}}{\partial T}\right)_V \Big|_{V_1(T)} + \left(\frac{\partial \bar{H}}{\partial V}\right)_T \Big|_{V_1(T)} \frac{dV_1}{dT} - \left(\frac{\partial \bar{H}}{\partial T}\right)_V \Big|_{V_0(T)} - \left(\frac{\partial \bar{H}}{\partial V}\right)_T \Big|_{V_0(T)} \frac{dV_0}{dT}$$

$$= -\frac{dP_S(T)}{dT}(V_1(T) - V_0(T)) - P_S(T) \left[\frac{dV_1(T)}{dT} - \frac{dV_0(T)}{dT} \right]$$

$-S(T, V_1(T))$

$-P_S(T) \frac{dV_1}{dT}$

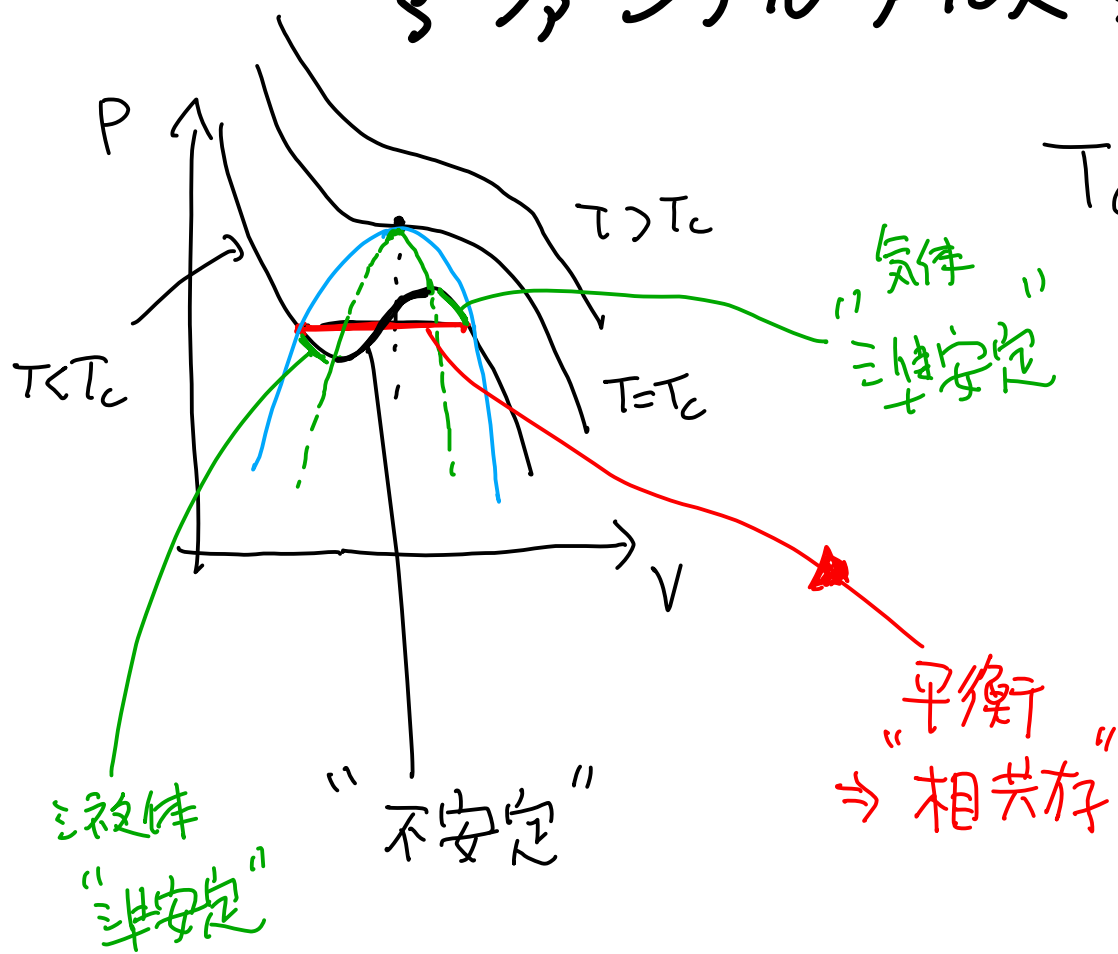


$$S(T, V_1(T)) - S(T, V_0(T)) = \frac{dP_S(T)}{dT}(V_1(T) - V_0(T))$$

$$\therefore \frac{dP_S(T)}{dT} = \frac{1}{T[V_1(T) - V_0(T)]}$$

~ Intermission ~

§ ファンデルワールス気体



$$T_c = \frac{8a}{27bR}$$

ファンデルワールス気体
状態方程式



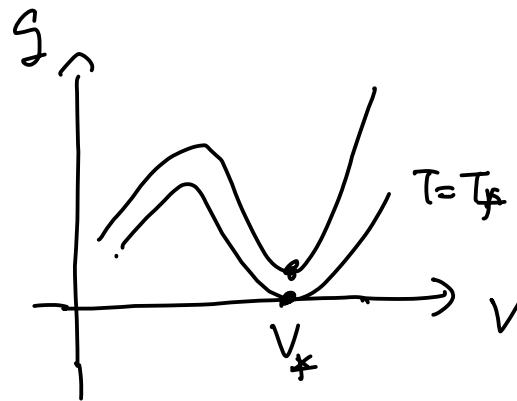
相共存?

- 熱力学変分原理



$$P = \frac{NRT}{V-Nb} - \left(\frac{N}{V}\right)^2 a \quad \S \text{計算}$$

$$\left(\frac{\partial P}{\partial V}\right)_T = -\frac{NRT}{(V-Nb)^2} + 2\frac{N^2}{V^3}a = 0$$



$$\begin{cases} g(V) \equiv NRT V^3 - 2N^2 a (V-Nb)^2 = 0 \\ g'(V) = 3NRT V^2 - 4N^2 a (V-Nb) = 0 \end{cases}$$

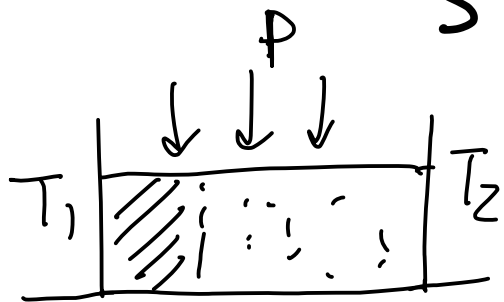
$$NRT V^3 - 2N^2 a (V-Nb)^2 - \frac{3}{2}NRT V^2 (V-Nb) + 2N^2 a (V-Nb)^2 = 0$$

$$V - \frac{3}{2}(V-Nb) = 0 \quad \underline{V_* = 3Nb}$$

$$T_* = \frac{2N^2 a}{NR V_*^3} (V_* - Nb)^2 = \frac{8N^2 a (Nb)^2}{27NR (Nb)^3}$$

$$= \frac{8a}{27Rb} //$$

§ 最近の話題



$$T_1 < T_0(p) < T_2$$

熱伝導下相共存

e.g. 水

$$T_1 = 95^\circ\text{C}$$

$$T_2 = 105^\circ\text{C}$$

P: 1気圧



界面の温度 $\approx 96^\circ\text{C}$

(理論的予言)

cf. Global thermodynamics
for heat conduction systems
J. Stat. Phys. 177 (2019)

§ 1.10-1

